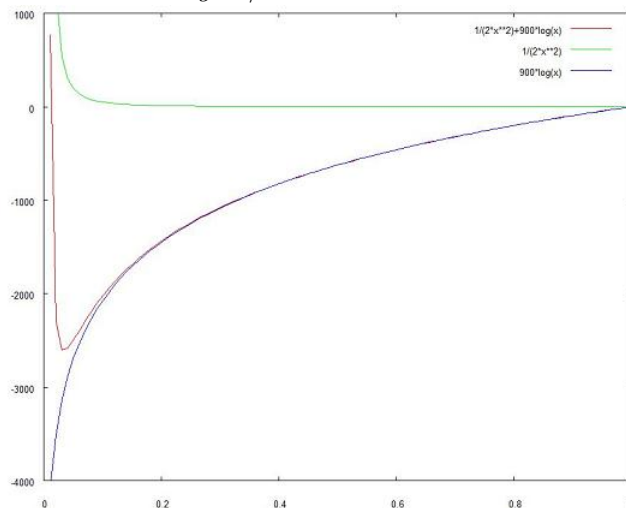


===== PROB 1 =====

$$V(r) = k \ln(r/L)$$

$$k = 900J; L = 1m$$

$$\text{angular momentum is } l = 1kgm^2/s$$



a) the plot

b) the minimum of the fictitious (effective) potential is when the motion is on a circular orbit.

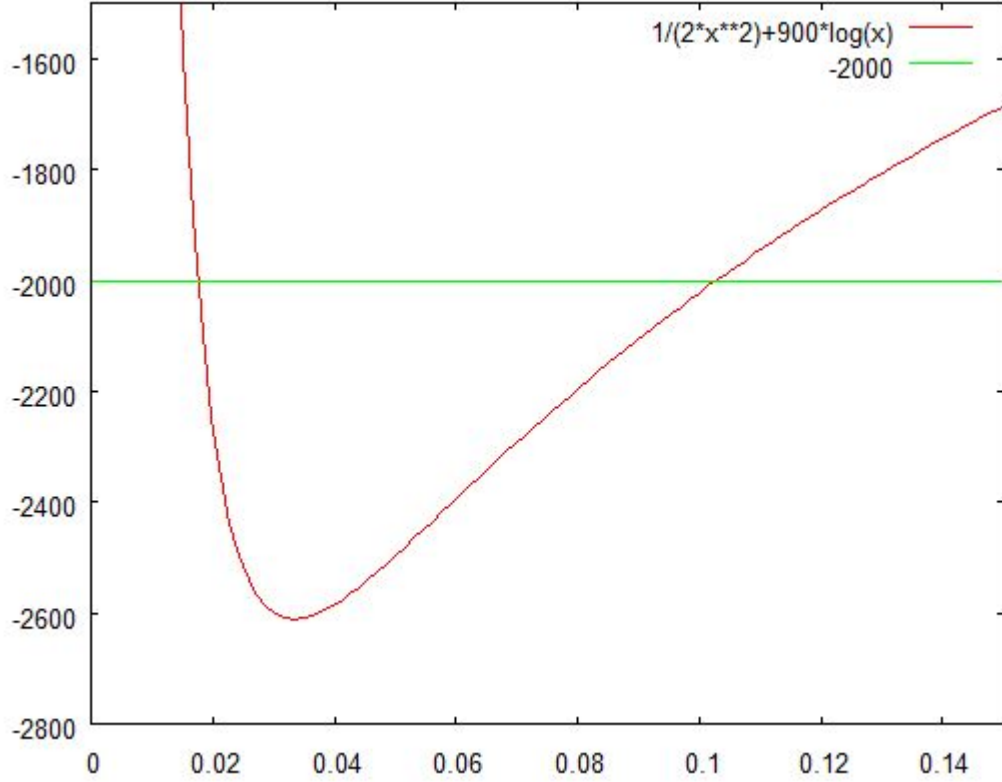
the fictitious potential is minimum when $-\frac{\partial V'(r)}{\partial r} = 0 = f(r) + \frac{l^2}{mr^3} = -\frac{\partial V(r)}{\partial r} + \frac{l^2}{mr^3} = 0$ see Goldstein p77

$$-\frac{\partial V(r)}{\partial r} = \frac{-k}{r} = \frac{-l^2}{mr^3}$$

assuming $l = 1kgm^2/s; k = 900J; m = 1kg$

$$r = \frac{l}{\sqrt{mk}} = \frac{1}{30}m$$

c) apsidal points are at the intersection of the energy with the fictitious potential



$E = V'(r) + \frac{1}{2}m\dot{r}^2 = k\ln(r/L) + \frac{1}{2}\frac{l^2}{mr^2} + \frac{1}{2}m\dot{r}^2$
 at the absidal points $E = V'(r)$ for it is $\dot{r}^2 = 0$
 hence by solving this eq $k\ln(r/L) + \frac{1}{2}\frac{l^2}{mr^2} = E$
 we find the wanted values of r.

$$r1 = 0.01743m$$

$$r2 = 0.10282m$$

$$\alpha = \frac{r1+r2}{2} = 0.06013$$

$$e = \frac{r2-r1}{r2+r1} = 0.71$$

===== PROB 2 =====

The energy of the planet before the impact is $E_i = \frac{1}{2}Mv_i^2 - \frac{k}{r}$ where v_i is the planet's initial speed, $r = r_{max}$ and $k = GM_\odot M$.

$E_f = \frac{1}{2}(M+m)v_f^2 - \frac{(M+m)k}{Mr}$ is the energy of the system (planet+comet) after the impact

Since this system is on parabolic orbit $E_f = 0 \implies v_f^2 = \frac{2k}{Mr}$ (1)

From the conservation of momentum we get $(M+m)v_f = Mv_i + mv_c$, where v_c is the speed of the comet before impact.

Hence $mv_c = (M+m)v_f - Mv_i$ (2).

Next we need to find v_i . Recall that the planet's orbit before impact is given by $r(\theta) = d/(1 + e \cos \theta)$, where the semi-latus rectum $d = \frac{l^2}{Mk}$. (3)

For $\theta = \pi$ we obtain the maxium distance $r = d/(1-e)$, and thus $d = (1-e)r$ (4).

Using (3) and (4) we see $l^2 = (1-e)rMk$ (5).

Yet l is also $l = rMv_i$, which we insert in (5) to obtain $M^2v_i^2 = (1-e)Mk/r = (1-e)M^2\frac{k}{Mr}$ (6).

Using (2), the kinetric energy of the comet before impact is $T_c = \frac{1}{2m}(mv_c)^2 = \frac{1}{2m}[(M+m)^2v_f^2 + M^2v_i^2 - 2(M+m)v_fMv_i] = \frac{1}{2m}[(M+m)^2\frac{2k}{Mr} + (1-e)M^2\frac{k}{Mr} - 2(M+m)M\sqrt{\frac{2k}{Mr}(1-e)\frac{k}{Mr}}] = \frac{1}{2m}[2(M+m)^2 + (1-e)M^2 - 2(M+m)M\sqrt{2(1-e)}]\frac{k}{Mr} = \frac{k}{2mMr}[2(M+m)^2 - 2\sqrt{2(1-e)}(M+m)M + (1-e)M^2]$

The comet's total energy before impact is $E_c = T_c - \frac{mk}{Mr} = T_c - \frac{k}{2mMr}2m^2$. Since $m \ll M$ we see that the second term is much smaller.

In fact if we let $m/M \rightarrow 0$ we find

$$E_c = T_c = \frac{k}{2mMr}M^2[2 - 2\sqrt{2(1-e)} + (1-e)] = \frac{kM}{2mr}[2 - 2\sqrt{2(1-e)} + (1-e)]$$

Using $r = (1+e)a$ and $e = 1 - \alpha$ we find

$$E_c = T_c = \frac{kM}{2m(2-\alpha)a}[2 - 2\sqrt{2\alpha} + \alpha].$$

Note: We had $T_c = \frac{kM}{2mr}[2 - 2\sqrt{2(1-e)} + (1-e)]$. If we set $r = d/(1-e)$, this becomes $T_c = \frac{kM}{2md}(1-e)(2 - 2\sqrt{2(1-e)} + (1-e))$. Since d is finite for $e \rightarrow 1$, we see that $T_c \rightarrow 0$ if $e \rightarrow 1$.

Note2: We could have gotten v_i also from $E_i = -\frac{k}{2a} = -(1+e)\frac{k}{2r}$ and the obvious $E_i = \frac{1}{2}Mv_i^2 - \frac{k}{r}$. We get again (6). Here the $E_i = -\frac{k}{2a}$ comes from adding the r_{min} and r_{max} found from the quadratic Eqn. $E_i = \frac{M}{2}(0 + r^2(\frac{l}{mr^2})^2) - \frac{k}{r}$ that gives us r_{min} and r_{max} .

===== PROB 3=====

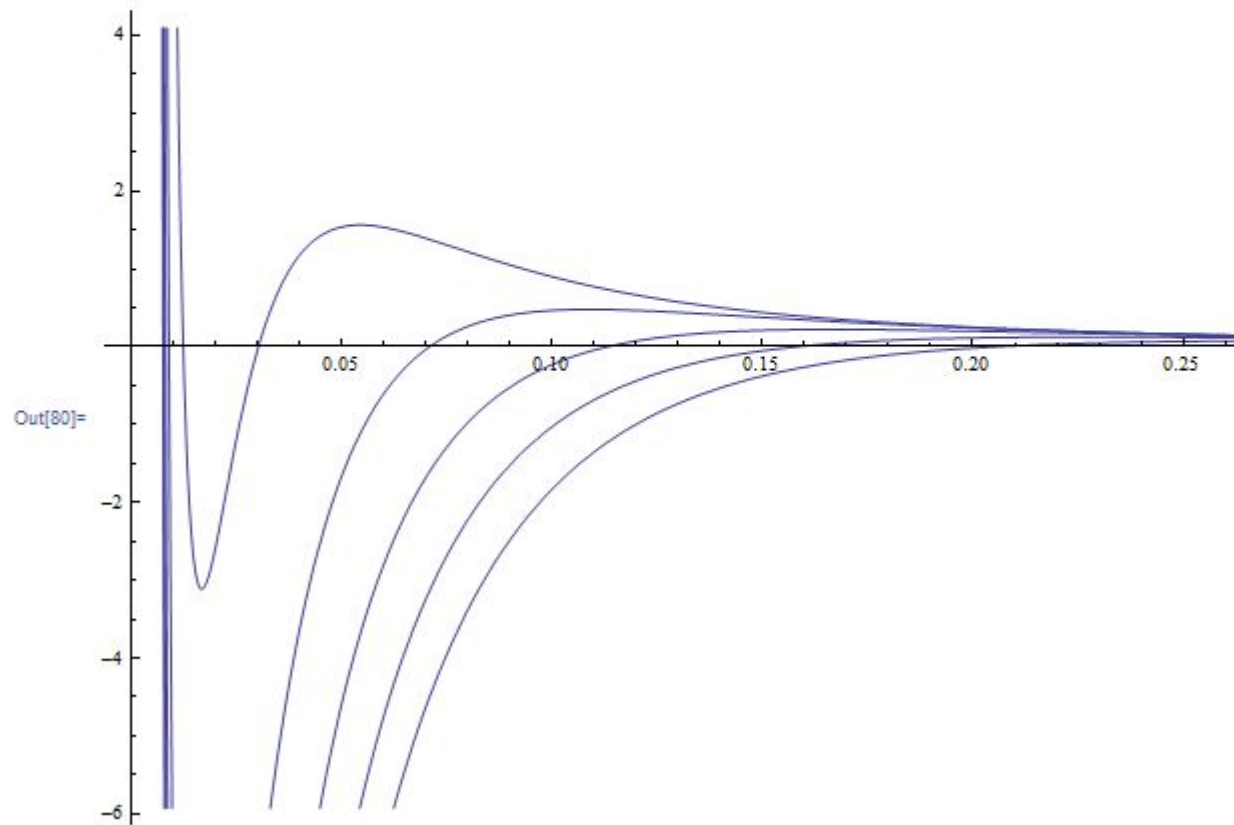
```

In[78]:= k = 1.5
         l = 0.1
         Plot[Table[ $\frac{l^2}{r^2} - \frac{k \star e^{-r/a}}{r}$ , {a, 0.02, 0.06, .01}], {r, 0, .3}]

```

Out[78]= 1.5

Out[79]= 0.1



The plot shows the effective potential for different values of a $0.02 \leq a \leq 0.06$.

Notice how at low values of $a = 0.02$ a maximum starts to show up.

The indicated values k, l, a, r are all fictitious.

In nuclear physics a is the range of the nuclear force: $\sim 10^{-15}m$

$$V(r) = -\frac{k}{r}e^{-\frac{r}{a}}$$

$$L = T - V$$

$$L = \frac{1}{2}m(\dot{r}^2 + r^2\dot{\theta}^2) + \frac{k}{r}e^{-\frac{r}{a}}$$

$$\frac{d}{dt}(\partial L_{\dot{\theta}}) - \partial L_{\theta} = 0$$

$$\partial L_{\dot{\theta}} = mr^2\dot{\theta}; \partial L_{\theta} = 0$$

hence $\frac{d}{dt}(\partial L_{\dot{\theta}}) = 0$

and $mr^2\dot{\theta} = \text{const}$

(1) $l = mr^2\dot{\theta}$ angular momentum of the system

$\frac{d}{dt}(\partial L_{\dot{r}}) - \partial L_r = 0$

$\partial L_{\dot{r}} = m\dot{r}; \partial L_r = mr\dot{\theta}^2 - \frac{k}{ar}e^{-\frac{r}{a}} - \frac{k}{r^2}e^{-\frac{r}{a}}$

$\frac{d}{dt}(\partial L_{\dot{r}}) = m\ddot{r}$

(2) $m\ddot{r} - mr\dot{\theta}^2 + \left(\frac{r}{a} + 1\right) \frac{k}{r^2}e^{-\frac{r}{a}} = 0$

(1) and (2) are the equations of motion

we can plug (1) into (2)

(3) $m\ddot{r} - \frac{l^2}{mr^3} + \left(\frac{r}{a} + 1\right) \frac{k}{r^2}e^{-\frac{r}{a}} = 0$

now that we know from (1) that l is a constant of the motion

so we can write the fictitious(effective) potential as

$V'(r) = \frac{l^2}{2mr^2} + V(r)$

$V'(r) = \frac{l^2}{2mr^2} - \frac{k}{r}e^{-\frac{r}{a}}$

$E = \frac{1}{2}m\dot{r}^2 + V'(r)$

$E = \frac{1}{2}m\dot{r}^2 + \frac{l^2}{2mr^2} - \frac{k}{r}e^{-\frac{r}{a}}$

----- prob 3part b)

use $u = 1/r$ and show that $\frac{d^2u}{d\theta^2} + u = F(u) = -(m/l^2) \frac{d[V(1/u)]}{du}$

$du/dr = -1/r^2$;

$\frac{du}{d\theta} = -\frac{1}{r^2} \frac{dr}{d\theta} = -\frac{1}{r^2} \frac{dr}{dt} \frac{dt}{d\theta} = -\frac{1}{r^2} \dot{r} \frac{1}{\dot{\theta}}$

hence from (1) $\frac{l}{mr^2} = \dot{\theta}$

$\frac{du}{d\theta} = -\frac{1}{r^2} \dot{r} \frac{mr^2}{l} = -\frac{m}{l} \dot{r}$

$\frac{d^2u}{d\theta^2} = \frac{dt}{d\theta} \frac{d}{dt} \left(-\frac{m}{l} \dot{r}\right) = -\frac{m}{l} \frac{1}{\dot{\theta}} \ddot{r} = -\frac{m}{l} \frac{m\ddot{r}}{l} \dot{r}$

$\frac{d^2u}{d\theta^2} = -\frac{m^2}{l^2} \ddot{r}$

$\ddot{r} = -\frac{d^2u}{d\theta^2} \frac{l^2}{m^2 r^2} = -\frac{d^2u}{d\theta^2} \frac{l^2}{m^2} u^2$ (6)

replacing (6) into (3) $m\ddot{r} - \frac{l^2}{mr^3} + \left(\frac{r}{a} + 1\right) \frac{k}{r^2}e^{-\frac{r}{a}} = 0$

$-\frac{d^2u}{d\theta^2} \frac{l^2}{m} u^2 - \frac{l^2}{m} u^3 + \frac{d}{dr} V(r) = 0$ where we replaced $\frac{d}{dr} V(r) = \left(\frac{r}{a} + 1\right) \frac{k}{r^2}e^{-\frac{r}{a}}$

$-\frac{d^2u}{d\theta^2} \frac{l^2}{m} u^2 - \frac{l^2}{m} u^3 - u^2 \frac{d}{du} V(u) = 0$

$\frac{d^2u}{d\theta^2} + u = -\frac{m}{l^2} \frac{d}{du} V(u) = F(u)$ which proves (i)

(ii) and (iii) Taylor expand.. $u = u_0 + \delta u$; $F(u_0) + F'(u_0)\delta u$

plug the above into the eq of motion $\frac{d^2u}{d\theta^2} + u = F(u)$ (7)

assume $u'' = \frac{d^2u}{d\theta^2}$; $F' = \frac{d}{du} F$

$u_0'' + \delta'' u + u_0 + \delta u = F(u_0) + F'(u_0)\delta u$

$u_0'' = 0$; for $u_0 = \text{const}$

from EOM (7) $u_0'' + u_0 = F(u_0) \Rightarrow 0 + u_0 = F(u_0)$ (8)

$\delta u'' + [1 - F'(u_0)]\delta u = F(u_0) - u_0 = 0$

$\delta u'' = -[1 - F'(u_0)]\delta u$

assuming $\beta^2 = [1 - F'(u_0)]$ (9)

$\frac{d^2u}{d\theta^2} = -\beta^2 u$

$\delta'' u = -\beta^2 \delta u$

the solution is $\delta u = A \cos \beta \delta \theta$

in order to find β we can use (i)

$$V(u) = -ku e^{\frac{-1}{au}}$$

$$F(u) = -\frac{m}{l^2} \frac{d}{du} V(u) = \left(-\frac{m}{l^2}\right) \left[-k - \left(\frac{1}{au^2}\right) \frac{ku}{a}\right] e^{\frac{-1}{au}}$$

$$F(u) = \frac{mk}{l^2} \left(1 + \frac{1}{au}\right) e^{\frac{-1}{au}} \quad (10)$$

now recall (8) $F(u_0) = u_0$

$$F(u_0) = \frac{mk}{l^2} \left(1 + \frac{1}{au_0}\right) e^{\frac{-1}{au_0}} = u_0$$

$$e^{\frac{-1}{au_0}} = u_0 \frac{l^2}{mk} \left(\frac{1}{1 + \frac{1}{au_0}}\right) \quad (11)$$

$$F'(u) = \frac{mk}{l^2} e^{\frac{-1}{au}} \left\{ \left[\frac{-1}{au^2}\right] + \left[1 + \frac{1}{au}\right] \left(\frac{1}{au^2}\right) \right\}$$

$$F'(u) = \frac{mk}{l^2} e^{\frac{-1}{au}} \left(\frac{1}{au^3}\right) = \frac{mk}{a^2 u^3 l^2} e^{\frac{-1}{au}} \quad (12)$$

$$F'(u_0) = \frac{mk}{a^2 u_0^3 l^2} e^{\frac{-1}{au_0}} \quad (13)$$

plug (11) into (13)

$$F'(u_0) = \frac{mk}{a u_0^3 l^2} \frac{u_0 l^2}{mk} \left(\frac{1}{1 + \frac{1}{au_0}}\right)$$

$$F'(u_0) = \frac{1}{a^2 u_0^2} \left(\frac{au_0}{1 + au_0}\right) = \frac{1}{au_0(1 + au_0)}$$

$$\beta^2 = [1 - F'(u_0)] = 1 - \frac{1}{au_0(1 + au_0)}$$

For $r_0/a \ll 1$ or $au_0 \gg 1$ we find $\beta^2 \approx 1 - \frac{1}{(au_0)^2} = 1 - (r_0/a)^2$ and thus

$$\beta \approx 1 - \frac{1}{2}(r_0/a)^2.$$

From $\beta \Delta\theta = 2\pi$ we obtain $\Delta\theta = 2\pi/\beta \approx 2\pi + \pi(r_0/a)^2$.

===== PROB 4 =====

given lunar month $T_{moon} = 27.3 \text{ days}$; earth period about the sun $T_e = 365.4$

days

mean radii of the earth's orbit about the sun $r_{earth} = 1.49 \times 10^{11} m$

mean radii of the moon's orbit about the earth $r_{moon} = 3.8 \times 10^8 m$

from Kepler III law the period of a planet is $T = 2\pi a^{3/2} \sqrt{\mu/k}$

assuming $\mu_{es} = \frac{m_e \times M_s}{m_e + M_s} \approx M_e$ the reduced mass of the sun, earth system

similarly assuming $\mu_{me} = \frac{m_m \times M_e}{m_m + M_e} \approx m_m$ the reduced mass of the earth,

moon system

$$k = GMm$$

a_{earth} earth's orbit semiaxis major about the sun

a_{moon} moon's orbit semiaxis major about the earth

$$T_e^2 = \frac{4\pi^2 a_e^3}{GM_s}, \text{ similarly } T_m^2 = \frac{4\pi^2 a_m^3}{GM_e}$$

$$\frac{T_m^2}{T_e^2} = \left(\frac{a_m}{a_e}\right)^3 \frac{M_s}{M_e}$$

$$\frac{M_s}{M_e} = \left(\frac{a_e}{a_m}\right)^3 \left(\frac{T_m}{T_e}\right)^2$$

plug in numbers to find that

$$\frac{M_s}{M_e} = \left(\frac{1.49 \times 10^8}{3.8 \times 10^5}\right)^3 \left(\frac{27.3}{365.4}\right)^2 = 3.365 \times 10^5$$

===== PROB 5 =====

$$f(r) = -\frac{\partial V}{\partial r}$$

$$\frac{1}{r^2} \frac{d}{d\theta} \left(\frac{l}{Mr^2} \frac{dr}{d\theta} \right) - \frac{l^2}{Mr^3} = f(r) \quad (\text{Goldstein 3.33})$$

$$\begin{aligned}
&\text{given } r = c\theta^2 \Rightarrow \frac{dr}{d\theta} = 2c\theta \\
&\frac{1}{r^2} \frac{d}{d\theta} \left(\frac{l}{Mr^2} 2c\theta \right) - \frac{l^2}{Mr^3} = f(r) \\
&\text{substitute } r^2 = c^2\theta^4 \\
&\frac{1}{r^2} \frac{d}{d\theta} \left(\frac{2lc\theta}{Mc^2\theta^4} \right) - \frac{l^2}{Mr^3} = f(r) \\
&\frac{1}{r^2} \frac{d}{d\theta} \left(\frac{-2l}{Mc\theta^3} \right) - \frac{l^2}{Mr^3} = f(r) \\
&\frac{1}{r^2} \left(\frac{-6l}{Mc\theta^4} \right) - \frac{l^2}{Mr^3} = f(r) \\
&\frac{1}{r^2} \left(\frac{-6l}{Mc\theta^4} \right) - \frac{l^2}{r^3} = f(r) \\
&\text{replace } r = c\theta^2 \\
&\frac{-6cl}{Mr^4} - \frac{l^2}{Mr^3} = f(r) \\
&V(r) = Ar^n + Br^m \\
&-\frac{\partial V}{\partial r} = -Anr^{n-1} - Bmr^{m-1} = f(r) = \frac{-6cl}{M}r^{-4} - \frac{l^2}{M}r^{-3} \\
&\text{hence } n = -3; m = -2
\end{aligned}$$